DOI: 10.1002/chem.200700067

## **Chemical Anisotropies of Carbon Nanotubes and Fullerenes Caused by the Curvature Directivity**

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**Abstract:** The directional-curvature theory is developed as a rational basis for the strain energy and the chemical reactivity in single-walled carbon nanotubes (SWCNTs) and fullerenes. The directional curvature  $K_{\rm D}$  and its mean  $K_{\rm M}$ , derived from this theory, cover the overall curvatures of their bonds and atoms and break through the limitations of the pyramidalized-angle  $\theta_{\rm p}$  approach, which is only available to

atomic curvature. The directional-curvature theory demonstrates that  $K_{\rm D}$  and  $K_{\rm M}$  depend directly on the strain or reactive binding energies of the bonds and atoms and that there is approximate curvature conservation in

**Keywords:** anisotropy • fullerenes • nanotubes • reaction mechanisms • reactivity • strain energy

SWCNTs and fullerenes. Application of this theory to addition reactions of various SWCNTs and fullerenes shows that the slope of the straight line between the strain or binding energies and  $K_D$  is close to a constant, which helps clarify the puzzle as to why some functionalizations of  $C_{70}$  occur at the relatively flat midsection.

### Introduction

Since fullerenes<sup>[1]</sup> and carbon nanotubes,<sup>[2]</sup> new types of carbon material different from graphite and diamond, were discovered, they have exhibited interesting electronic, electric, mechanical, and structural properties and stimulated extensive research into their chemical functionalization.<sup>[3]</sup> Experimental and theoretical investigations show that fullerenes and single-walled carbon nanotubes (SWCNTs) have higher curvature-induced chemical reactivities than flat graphite<sup>[4]</sup> and have a different chemical anisotropy from graphite and diamond.<sup>[5]</sup> Research into the strain or curvature, which is considered to be the source of their higher reactivity, has become an important aspect of the chemistry of SWCNTs and fullerenes. Theoretical research into the corre-

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lation of curvature with the reactive mechanisms, up until now, remains at the level of isolated atoms<sup>[4b-c,5a]</sup> and is deficient in various directional curvatures, hence the reactivity of the bond and the inherent relation of the curvature of the atom and bond cannot be considered correctly. Consideration of the whole curvature of these carbon clusters is indispensable to rationally estimate the reactivities of the atoms and bonds.

The formation of SWCNTs or fullerenes with various frameworks includes two orbital-interaction processes. First, the orbitals of isolated carbon atoms are hybridized to sp<sup>x</sup> (type:  $\sigma$ ; s/p orbital ratio: 1:x) and sp<sup>y</sup> (type:  $\pi$ ; s/p orbital ratio: 1:y). Second, the multifarious carbon tubes or cages are formed by different routes between the hybrid orbitals of the adjacent atoms. The second precedence is the crux of the formation and anisotropic reactivity of the SWCNTs or fullerenes with a specific structure (Figure 1).

The pyramidized-angle  $(\theta_p)$  concept<sup>[4b,6]</sup> indicates the orbital hybridization of isolated atoms; however, it only considers the first step, so that the second step, the formation of SWCNTs and fullerenes, exceeds the capacity of this approach. The misalignment angle based on pyramidization also does not represent the second step correctly. Therefore, the pyramidization approach to carbon atoms merely represents the mean reactivity of the vertex atom of the fullerene or SWCNT. This approach has undergone grave difficulty in describing the reactivities and adduct structures of the different directional C–C bonds on SWCNTs.<sup>[5a–b]</sup> Furthermore,







Figure 1. SWCNT, fullerene, graphene, and diamond formed by various routes.

the pyramidization approach also does not have the proper criterion for the reactivities of fullerenes with anisotropic curvature, which encompasses most of the known fullerenes. For example, the tubular radii and  $\theta_{\rm p}$  value of SWCNTs (5,5) and (9,0) are very close, whereas the curvatures and reactivities of the C-C bonds on the circumference of the (5,5) tube are much greater than those of the axis bonds of the (9,0) tube because of the different routes of formation.<sup>[5b,7]</sup> The reactivities of the variously oriented C-C bonds in the (5,5) or (9,0) tubes are also obviously different as a result of the different oblique angle  $\theta^{[5b]}$  of the C-C bonds, although their  $\theta_p$  values are equal. Diamond and  $C_{20}$ also have similar  $\theta_p$  values, whereas their structures are quite distinct. From the addition reactions of benzyne, dichlorocarbene, and sulfonyl azides to  $C_{70}$ <sup>[5c-e]</sup> it has been found that the sites on  $C_4-C'_4$  (i.e.,  $C_7-C_8$  in IUPAC numbering) with smaller pyramidalized angles have higher reactivities than those with larger  $\theta_p$  values. The [2+1] addition of an oxygen atom to C50 does not occur at the equator with the largest  $\theta_{\rm p}$  value.<sup>[8]</sup>

The bond-curvature approach<sup>[5b]</sup> was successfully used to explain the reactivities of the C–C bonds of SWCNTs; however, this approach needs to be extended and enriched to predict the reactivities of the fullerenes and vertex sites of SWCNTs. To describe the oriented curvature and the anisotropic chemistry of SWCNTs and fullerenes generated by various routes of formation, we have developed, herein, a unified, directional-curvature theory that evaluates the strain energies and reactivities of the vertices and bonds.

# **FULL PAPER**

#### Directional-Curvature Theory of SWCNTs and Fullerenes

The direction of the curvature of the  $\sigma$  framework camber of SWCNTs or fullerenes is in the normal direction of the camber, the orientation of which is the  $\pi$  hybrid orbital relative to the σ framework formed by the delocalized interaction of the hybrid orbitals of the isolated atoms. Hence, the curvature of the camber can be regarded as a unified mathematical representation of the geometry and inherent stress of the dangling hybrid orbital on the vertex or the arc of the camber relative to the  $\sigma$  framework of these carbon clusters as viewed by delocalization.

The curvatures of the arcs that correspond to the C–C bonds and the vertex carbon

atoms of the SWCNTs or fullerenes are defined as the directional curvature  $K_{\rm D}$  and its mean  $K_{\rm M}$ , respectively. Based on the delocalization of the curved  $\pi$  conjugated system<sup>[5b]</sup> and the Euler formula in differential geometry<sup>[9]</sup> for the C<sub>a</sub>-C<sub>b</sub> bonds with biterminal oblique angles  $\theta_{\rm i}$ ,<sup>[5b]</sup>  $K_{\rm D}$  can be defined by Equation (1) and for the vertex atom i, the curvature may take the mean of the  $K_{\rm D}({\rm i})$  for  $\theta_{\rm i}$ , namely,  $K_{\rm M}$ , which is given by Equation (2):

$$K_{\rm D} = \frac{1}{2} \sum_{i=a,b} \left( K_{i1} \sin^2 \theta_i + K_{i2} \cos^2 \theta_i \right)$$
(1)

$$K_{\rm M} = \frac{2}{\pi} \int _{0^{\frac{\pi}{2}}} K_{\rm D}(\mathbf{i}) d\theta_{\mathbf{i}} = \frac{1}{2} (K_1 + K_2)$$
(2)

where  $K_{i1}$  and  $K_{i2}$  and  $K_1$  and  $K_2$  are the principle curvatures of carbon atom i (see the Supporting Information). Equations (1) and (2) include the whole curvature of the vertices and the bonds of the various SWCNTs or fullerenes. In practice, others may be simplified greatly except for the curvature of the ellipsoidal fullerenes, which should be calculated by Equations (1) and (2). For example, in the SWCNTs, because  $K_1$  is equivalent to 1/R (see the Supporting Information) and  $K_2=0$  and  $\theta_a = \theta_b$ ,  $K_D$  is defined by Equation (3):

$$K_{\rm D} = \frac{\sin^2 \theta}{R} \tag{3}$$

where *R* is the radius of the SWCNT. Equation (3) is the bond curvature *K* of the SWCNTs,<sup>[5b]</sup> and the  $K_{\rm M}$  value for

Chem. Eur. J. 2007, 13, 6430-6436

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Table 1. The curvature parameters of SWCNTs in the formation of the adducts with the opened C-C bond.

SWCNTs	Atoms						Bonds					
	R	$K_{\rm M}$	$\theta_{\rm p}$	$\sum K_{\rm M}$	Reactivity	$\theta$	$K_{\rm D}$		$E_{\rm b}  [ m kcal  mol^{-1}]$		Reactivity	
	[nm]	$[nm^{-1}]$	[°]	$[n^{j}m^{-1}]$		[°]	$[nm^{-1}]$	O/SWCNT	CH <sub>2</sub> /SWCNT <sup>[a]</sup>	NH/SWCNT <sup>[a]</sup>		
(11,0)	0.431	1.16	4.70	22.10		60.00	1.742	71.89				
(6,2)	0.282	1.77	7.17	22.07		46.10	1.841	73.20	68.38	42.71		
(8,8)	0.542	0.92	3.73	22.13		90.00	1.844	78.86				
(9,0)	0.352	1.42	5.75	22.17		60.00	2.129	79.43				
(8,0)	0.313	1.60	6.48	22.16		60.00	2.396	82.14			weak	
(6,1)	0.257	1.95	7.90	22.14	$\approx K_{ m M}$	52.41	2.429	85.48				
(6,6)	0.407	1.23	4.98	22.17	or	90.00	2.457	89.42				
(7,0)	0.274	1.82	7.40	22.20	1	60.00	2.737	86.81				
(6,4)	0.341	1.47	5.95	22.05	2R	83.41	2.894	93.27			•	
(5,5)	0.339	1.47	5.95	22.14		90.00	2.950	96.79			strong	
(6,3)	0.311	1.61	6.52	22.13		79.11	3.101	96.98	91.86	65.96		
(6,0)	0.235	2.13	8.63	22.07		60.00	3.191	95.12				
(6,2)	0.282	1.77	7.17	22.16		73.90	3.273	97.81	94.62	68.27		
(6,1)	0.257	1.95	7.90	22.20		67.59	3.325	100.33				
(4,4)	0.271	1.85	7.49	22.20		90.00	3.690	104.50				

[a] The binding energies are given on the basis of the cluster model.<sup>[5b]</sup>

the vertices of SWCNTs can be obtained from Equation (2) or the average of the  $K_{\rm D}$  value with  $\theta_{\rm min}$  and  $60^{\circ} \pm \theta_{\rm min}$ ,<sup>[5b]</sup> namely from Equation (4)

$$K_{\rm M} = \frac{1}{2R} \tag{4}$$

hence, the vertex curvatures of SWCNTs with a given radius are identical everywhere. For the spherical fullerenes, such as  $C_{60}$ ,  $K_D$  and  $K_M$  are given by Equation (5):

$$K_{\rm D} = K_{\rm M} = \frac{1}{R_{\rm S}} \tag{5}$$

because  $K_1 = K_2 = 1/R_s$  (see the Supporting Information), where  $R_s$  is the radius of spherical fullerenes.

It can be seen, from the above discussion, that the curvatures of various directions are different from those of the isolated atoms in SWCNTs and that for the great majority of fullerenes  $K_{\rm D} \neq K_{\rm M}$ , except for the spherical fullerenes. In other words, as the SWCNTs or fullerenes are formed, the  $\sigma$ and  $\pi$  hybrid orbitals of a central carbon atom (set as C<sub>1</sub>) may be changed from the original sp<sup>x</sup> ( $\sigma$ ) and sp<sup>y</sup>( $\pi$ ) hybrid orbitals obtained by POAV analysis  $^{[6]}$  to new  $\sigma$  and  $\pi$  hybrid orbitals as a result of various curvatures of the  $C_1-C_n$  (n=2, 3, and 4) bonds between  $C_1$  and its adjacent atoms  $C_2$ ,  $C_3$ , or C<sub>4</sub>. At the same time, the  $\theta_{o\pi}$  value (or  $\theta_p + \pi/2$ ) of the C<sub>1</sub> atom will be changed correspondingly. For example, for an armchair SWCNT with a moderate radius, the  $\theta_{\sigma\pi}$  value for  $C_1$  obtained from the  $\sigma$  bond angles  $\gtrless C_n - C_1 - C_n'$ , obtained on the basis of the POAV approach,<sup>[6]</sup> comes close to equal. However, the practical  $\theta_{\alpha\pi}$  value (set as  $\theta_{\alpha\pi'}$ ) may be obviously different because of the curvature directivity. The  $\theta_{p}'$  value (or  $\theta_{\sigma\pi'} - \pi/2$ ) of C<sub>1</sub> along the C<sub>1</sub>-C<sub>n</sub> bond (set as  $C_1-C_2$ ) of the circumference is twice as large as the original  $\theta_{p}$  value, and the  $\theta_{p'}$  value along the C<sub>1</sub>-C<sub>3</sub> or C<sub>1</sub>-C<sub>4</sub> bonds

is one half of the  $\theta_p$  value. These  $\theta_{\sigma\pi'}$  values will create a set of new  $\sigma$  and  $\pi$  hybrid orbitals of the C<sub>1</sub> atom under orthonormal conditions.

The variance of the  $\theta_{o\pi}$  value caused by the curvature or strain led to increasing ( $\theta_{o\pi}$  value decreased) or decreasing ( $\theta_{\sigma\pi}$  value increased) bonding interactions between the  $\pi$  hybrid orbitals of the C<sub>1</sub> and C<sub>n</sub> atoms, and at the same time, the  $\sigma$  bond between C<sub>1</sub> and C<sub>n</sub> is changed reasonably. From the point of view of energy, the SWCNTs or fullerenes will tend to stabilize, but there is a larger difference in the activities of various C<sub>1</sub>-C<sub>n</sub> bonds; for example, the reactivity of the circumference C<sub>1</sub>-C<sub>2</sub> bond is obviously stronger than that of the axis C<sub>1</sub>-C<sub>3</sub> or C<sub>1</sub>-C<sub>4</sub> bonds. From the correlation of the directional curvature with the calculated binding energies and experimental results (see Table 1, Table 2, and reference [5]), it must be admitted that the directional-curvature method is reasonable, though the direct experimental measurement of the  $\theta_{\sigma\pi}'$  or  $\theta_{\sigma\pi}$  values is difficult.

The hybrid-orbital recombination inevitably leads to the anisotropy of the chemistry; however, the chemistry of the graphene and diamond is isotropic, because the hybrid orbitals are not changed. This behavior is just the chemical outcome of the directional-curvature theory.

#### The Directional-Strain Energy Derived from the Directional Curvature

The recombination of the atomic hybrid orbitals in SWCNTs and fullerenes is carried out according to their curvatures, which can be used to represent the interactional energy of these hybrid orbitals relative to a single sheet of graphite, as the SWCNT or fullerene increases the bonding energy and, at the same time, bears a corresponding strain energy that comes from the curved structure. Based on a continual elastic model,<sup>[10]</sup> the strain energy or work  $E_s$  re-

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Table 2. The curvature parameters for oxygen addition to fullereness

Fullerenes	Atoms						Bonds							
	Atomic	$R_{\rm s}$	$K_{\rm M}$	$\theta_{\rm p}^{\rm e} (\theta_{\rm p}^{\rm s})$	$\sum K_{\rm M}^{2}$	Reactivity	Bond	$K_{\rm D}$	$\theta_{p}^{e}(B)$	r	$E_{\mathrm{b}}$	Reactivity		
	types	[nm]	$[nm^{-1}]$	[°]	$[\dot{n}m^{-1}]$		types	$[nm^{-1}]$	[°]	[nm]	[kcal mol <sup>-1</sup> ]	opened C–C bond	closed C-C bond	
C <sub>60</sub>	1	0.352	2.84	11.50	483.9		[6,6] <sup>[a]</sup>	2.84		0.138	76.33		[b]	
							[5,6] <sup>[a]</sup>	2.84		0.145	78.25	[b]		
	1		3.109	12.59			$1(C_1 - C'_1)[5,6]$	3.141	9.72	0.1452	78.18	[b]		
	2		2.804	11.36		strong	$2(C_1 - C_2) [6,6]$	2.805		0.1397	75.89		[b]	
$C_{70}$	3		2.601	10.53	484.8	<b>Λ</b>	$3(C_2 - C_3)[5,6]$	2.779	9.45	0.1448	75.74			
							$4(C_3 - C'_3)$ [6,6]	2.899		0.1389	75.30		[b]	
	4		2.426	9.83			$5(C_3 - C_4)$ [5,6]	2.236	8.80	0.1450	70.43			
						weak	6(C <sub>4</sub> -C' <sub>4</sub> ) [5,6]	2.799	8.11	0.1434	76.97	[b]		
	5		2.375	9.62			$7(C_4 - C_5)[6,6]$	2.226		0.1421	63.46			
							$8(C_5 - C'_5)[6,6]$	2.769		0.1471	89.02		[b]	
	1		2.538	10.28		weak	$1(C_1 - C'_1)[5,6]$	2.520		0.1465	80.11			
	2		2.892	11.71			$2(C_1 - C_2) [6,6]$	2.861		0.1390	81.76		[b]	
C <sub>50</sub>	3		3.256	13.19	485.0	$\downarrow$	$3(C_2 - C_3)$ [5,6]	2.964		0.1469	88.31	[b]		
							$4(C_{2}-C_{2})$ [6.6]	2.816		0.1415	69.58		[b]	
	4		3.535	14.32		strong	$5(C_2-C_4)$ [5.6]	3.669		0.1417	82.80	[b]		
						8	$6(C_4 - C'_4)$ [5.5]	2.913		0.1470	73.48			

[a] [6,6] or [5,6] ring fusion. [b] Reaction is favorable.

quired to bend a bar with moment M through an angle  $\phi$  is given by Equation (6):

$$E_{\rm s} = \frac{1}{2}M \times \phi = \alpha \frac{\sin^2 \theta}{R} = \alpha K_{\rm D} \tag{6}$$

where  $\phi = \frac{r_{C^-C}}{R} \sin^2 \theta$ ,<sup>[5b]</sup> or more generally,  $\phi = r_{C^-C} K_D$ , in which  $r_{C^-C}$  is the corresponding arc of the C–C bond, and  $\alpha$  may be taken approximately as a constant because *M* does not significantly affect the  $\pi$  orbital, which is the dominant contributor to  $E_b$ . Thus, the mean  $E_s$  ( $E_{sM}$ ) is equal to  $\alpha K_M$ .

Equation (6) develops a relation in which the strain energy depends directly on  $K_{\rm D}$  or  $K_{\rm M}$ . The directional-curvature theory allows the curvature to represent the strain energy or the difference of the binding energy  $E_{\rm b}$  of the SWCNT or fullerene relative to the graphite sheet  $E_{\rm bg}$ , as given by Equation (7) or Equation (8):

$$E_{\rm b} = E_{\rm bg} + E_{\rm s} = E_{\rm bg} + \alpha K_{\rm D} \tag{7}$$

$$E_{\rm b} = E_{\rm T}({\rm C}_n) + E_{\rm X} - E_{\rm T}({\rm C}_n + {\rm X}) \approx E_{\rm T}({\rm C}_n') + E_{\rm C} + E_{\rm s} + E_{\rm X}$$
$$-(E_{\rm T}({\rm C}_n' + E_{\rm C-X} + E_{\rm X}) = (E_{\rm C} - E_{\rm C}' - E_{\rm C-X}) + E_{\rm X} = E_{\rm bg} + E_{\rm s}$$
(8)

where  $E_{\rm T}(C_n)$  is the energy of a fictitious SWCNT or fullerene from which the carbon atom or C–C bond participating in the reaction has been removed;  $E_{\rm C}$  and  $E'_{\rm C}$  are the energies corresponding to the carbon atom or bond before and after the reaction; and  $E_{\rm X}$  and  $E_{\rm C-X}$  are the energies of applied species x and the binding energy of the bond C–X, respectively. In the sidewall additions of H and Al atoms on the vertices of zigzag<sup>[11]</sup> and armchair<sup>[12]</sup> SWCNTs and of O, CH<sub>2</sub>, NH, and SiH<sub>2</sub> species on the C–C bonds with different  $\theta$  values for the various types of SWCNTs,<sup>[5b]</sup> the linear relationships between  $E_{\rm b}$  and 1/R,  $\theta_{\rm p}$ , or  $\frac{\sin^2 \theta}{R}$  have been seen, respectively.

#### The Relationship of $K_{\rm D}$ with the Bond Curvature K and the Pyramidalized Angle $\theta_{\rm p}$

From Equations (1), (3), and (5) deduced from Equation (1), it can be seen that  $K_D$  is the same as the bond curvature Kand its extension to the SWCNTs and fullerenes. It can be also proved that the  $K_M$  values from Equations (2), (4), and (5) represent the curvatures that correspond to the  $\theta_p$  value of the vertex atoms of SWCNTs and fullerenes, respectively.

Figure 2 is the projection of the bonds that adjoin a vertex carbon atom on the circumference of an armchair SWCNT. The  $\Delta C_1$ - $C_2$ - $C_3$  of the bottom plane is close to an equilateral triangle, thus giving Equation (9) and Equation (10)

$$K_{\rm M} = \frac{1}{2R} = \frac{\sin\frac{\phi}{2}}{r} \approx \frac{\phi}{2r} = \frac{2\theta_{\rm p}}{r} \tag{9}$$

$$K_{\rm D} = 2K_{\rm M}\sin^2\theta = \frac{4\theta_{\rm p}}{r}\sin^2\theta \tag{10}$$

where *r* is the C–C bond length of about 0.142 nm.

Chem. Eur. J. 2007, 13, 6430-6436

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Figure 2. The projected frame of the bonds on the circumference of an armchair SWCNT.

In Equation (9),  $\theta_p = \phi/4$  and, simultaneously, may be regarded as the mean of the radial and axial  $\theta_{p}$  value, in which the former is  $\phi/2$  and the latter zero as a result of orbital delocalization. Furthermore,  $\theta_p = \frac{r}{4R} = \frac{a}{4\sqrt{3R}}$  (a is the length of the repeat unit) can be also obtained, which is identical to the result reported by Kleiner and Eggert.<sup>[13]</sup> The  $K_{\rm M}$  and  $\theta_{\rm p}$  values are only dependent on the tubular radius R, and the  $\theta_p$  values obtained from  $K_M$  are also consistent with the result of Haddon and co-workers,<sup>[5a]</sup> but the tedious calculation is not required; besides, the vertex curvatures obtained from the  $K_M$  and  $\theta_p$  values are identical as well. Equations (9) and (10) may be spread to other SWCNTs or fullerenes if only the circular section in Figure 2 is taken as the corresponding curvature circle. Therefore, the relationship of  $K_{\rm D}$  (or  $K_{\rm M}$ ) with  $\theta_{\rm p}$  in the zigzag or chiral SWCNTs can be also represented by Equations (9) and (10), whereas those of the spherical fullerenes are given by Equation (11):

$$K_{\rm M} = K_{\rm D} = \frac{2\theta_{\rm p}^{\rm s}}{r} \tag{11}$$

where  $\theta_p^s$  is the pyramidalized angle of the vertices in the spherical fullerenes. For the ellipsoid fullerenes, those are given by Equation (12) and (13):

$$K_{\rm M} = \frac{1}{r} (\theta_{\rm p1}^{\ \rm e} + \theta_{\rm p2}^{\ \rm e}) = \frac{2\theta_{\rm p}^{\ \rm e}}{r}$$
(12)

$$K_{\rm D} = \frac{1}{r} \sum_{\rm i=a,b} \left( \theta_{\rm p1}^{\rm e} \sin^2 \theta_{\rm i} + \theta_{\rm p2}^{\rm e} \cos^2 \theta_{\rm i} \right)$$
(13)

where  $\theta_p^{e}$ ,  $\theta_{p1}^{e}$ , and  $\theta_{p2}^{e}$  are the corresponding pyramidalized angles of the vertex,  $K_1$  and  $K_2$ , on the ellipsoid fullerenes, respectively. It is evident from Equations (9)–(13) that the  $\theta_p$ ,  $\theta_p^{s}$ , or  $\theta_p^{e}$  values may represent the  $K_M$  value of the vertices in SWCNTs and fullerenes, although they are not enough to represent the directional curvatures  $K_D$  (see Table 1), except for the spherical fullerenes, and that the  $\theta_p^{s}$  value is equal to twofold  $\theta_p$ , as the *R* values are equal. For example, as a result of the close *R* values,  $\theta_p \approx 6^{\circ}$  for the (5,5), (6,4), and (9,0) tubes and  $\theta_p^{s} \approx 12^{\circ}$  for  $C_{60}$ . It can be found from Equation (9) that the sum of the curvatures of the vertices in the SWCNTs with i atoms is given by Equation (14)

$$\sum_{i} K_{\rm M} = m \sum_{j} K_{\rm M} = m \sum_{j} \frac{\phi}{2r} = m \frac{\pi}{r} = 22.13 \, m (nm^{-1})$$
(14)

where *j* and *m* are the atomic number in the circumference camber fragment with continuous C–C bonding and the repeated number of the camber fragment along the tubular axis, respectively. Therefore, there is also an approximate curvature conservation in SWCNTs as *m* is an infinite or the same certain value (see Table 1 and the Supporting Information). The curvature conservation in fullerenes may be obtained by an approach used by Haddon<sup>[4b]</sup> and  $\sum K_M^2$  (or  $\sum \theta_p^2$ ; see Table 2). Unfortunately, the curvature <sup>*i*</sup> conservation of SWCNTs cannot be represented by this approach.

#### Theoretical Calculations of the Directional Curvature and Prediction of Reactivity

The directional-curvature theory allows a universal method that estimates the reactivities of the vertices and bonds of SWCNTs and fullerenes to be developed (see Tables 1 and 2). The  $K_{\rm D}$  and  $K_{\rm M}$  values given in Tables 1 and 2 are calculated from the above-mentioned equations and the differential-geometry approach based on optimized structures. For SWCNTs and fullerenes, periodic<sup>[14]</sup> and molecular<sup>[15]</sup> DFT methods were applied to optimize the structures and calculate the binding energies, respectively (see the Supporting Information for details).

As an illustration of the directional-curvature theory, the calculated results for various SWCNTs with moderate R values show that there is a linear relationship between the  $E_{\rm b}$  and  $E_{\rm s}$  or  $K_{\rm D}$  values; for example,  $E_{\rm b}=e_{\rm bg}+\alpha K_{\rm D}=45.45+16.26 K_{\rm D}$  for the addition of oxygen (Figure 3). The intercept of the line could be regarded as the binding energy of the addition of oxygen onto the graphene plane; hence, the strain energy  $E_{\rm s}=\alpha K_{\rm D}$  in SWCNTs accounts for about 100% of the excess energy of the cured structure relative to graphene.<sup>[11]</sup>

Similar results can be obtained from different computational models and methods.<sup>[5b]</sup> It is noteworthy that the slopes  $\alpha$  are very similar for the various models, computational approaches, or adducts. For example, from the cluster model and DFT-B3LYP method<sup>[5b]</sup>  $\alpha = 16.96$  kcal nm mol<sup>-1</sup> for the addition of oxygen and  $\alpha = 16-18$  kcal nm mol<sup>-1</sup> for various adducts O, CH<sub>2</sub>, or NH with triplet ground states. The  $\alpha$  values are also equal for the vertex additions of H and Al atoms.<sup>[11]</sup> It can also be seen from Tables 1 and 2 that the conservation of  $\sum K_{\rm M}$  in various types of SWCNTs or of  $\sum K_{\rm M}^2$  in various types of fullerenes occurs, and that the latter<sup>1</sup> is approximately equal to the square of the former. In

#### 120 114 108 102 $E_{\rm h}$ (kcal mol<sup>-1</sup>) 96 90 84 $E_{\rm b} = 45.45 + 16.26 K_{\rm D}$ Correlation coefficient: 97.4% 78 72 66 60 1.5 2.0 3.0 3.5 4.0 2.5 $K_{\rm D}$

Figure 3. Variation of the binding energy  $(E_b)$  as a function of the  $K_D$  value of oxygen adducts with an opened C–C bond on SWCNTs.

addition, the  $\theta_p$  (or  $\theta_p^s$ ,  $\theta_p^e$ ) values obtained from  $K_M$  and Niyogi et al.<sup>[5a]</sup> are close.

As distinct from SWCNTs, the C–C bonds in fullerenes may result from either [6,6] or [5,6] ring fusion (see the Supporting Information and Table 2 for oxygen addition to fullerenes) and the change in the bond length is larger. In general, the  $E_{\rm b}$  value of the bonds with the [6,6] ring fusion or the shorter bond lengths are less because larger energies are required to break these bonds. The anisotropy of the chemistry of C<sub>60</sub> may be expressed by the different reactivities of the various bonds of the [6,6] and [5,6] ring fusions, the vertices, or the bonds of the distorted C<sub>60</sub>. The bonds of the [5,6] ring fusion are more reactive than those of the [6,6] ring fusion.<sup>[16]</sup> This case can be also seen from the calculated results for the defect SWCNT.<sup>[17]</sup>

The calculated results of the  $K_{\rm D}$  values for  $C_{70}$  (see Table 2 or the Supporting Information) rationally explain the addition<sup>[5e]</sup> of sulfonyl azides to  $C_{70}$  that occur at the  $C_1$ - $C_1'$  and  $C_4$ - $C_4'$  bonds (i.e.,  $C_2$ - $C_3$  and  $C_7$ - $C_8$  according to the IUPAC numbering). These bonds ( $C_1$ - $C_1'$  and  $C_4$ - $C_4'$ ) with [5,6] ring fusion have the largest  $K_{\rm D}$  values or binding energies for the formation of the fulleroids with an open C-C bond structure.

Up until now, no one has succeeded in explaining the puzzle of the higher reactivity of the  $C_4-C_4'$  bond at the relatively flat midsection, our approach based on the theory of  $K_D$  contributes to clarifying this dilemma. In fact, the  $K_D$  value is not simply equal to the average of the biterminal atomic  $K_M$  value of the bond, and it is also not equal to the curvature obtained by the average  $\theta_p^{e}(B)$  of the  $\theta_p^{e}$  value relative to the bonds in the POAV2 analysis<sup>[6]</sup> (see Table 2), thus the  $K_D$  value may be larger, even if the biterminal  $K_M$ ,  $\theta_p^{e}$ , or  $\theta_p^{e}(B)$  values are less. This reason is why the  $\theta_p^{e}$  value could not correctly describe the reactivities of the C-C bonds with various orientations and sites.

The  $K_{\rm D}$  value also indicates that the adducts of  $C_{70}$  oxidoannulenes with an open C–C bond structure formed at the  $C_1$ – $C_1'$  and  $C_4$ – $C_4'$  bonds are favorable.<sup>[18]</sup> The especially large  $E_{\rm b}$  value of the equatorial  $C_5$ – $C_5'$  bond may relate to

the largest bond length, and the absence of adduct  $C_5, C_5' - C_{70}O^{[19]}$  may be attributed to kinetic factors as the electronic state of the equatorial carbon atoms is very small at the HOMO/LUMO (see the Supporting Information).

There is also an approximate linear relationship between the  $E_{\rm b}$  and  $K_{\rm D}$  values for the bonds of the [5,6] ring fusion with the close C–C bond length in C<sub>70</sub> (Figure 4). For the



Figure 4. Variation of the binding energy  $(E_b)$  as a function of the  $K_D$  value for the bonds with [5,6] ring fusion during the addition of an oxygen atom.

vertex reactivity of  $C_{50}$ , the results predicted from the  $K_M$ and  $\theta_p^{e}$  values are the same, hence the [1+1] addition of a Cl atom onto  $C_{50}$  occurs preferably at C<sub>4</sub>, which is consistent with the experimental results.<sup>[19]</sup> However, the reactivities predicted by the  $K_D$  and  $\theta_p$  (or  $K_M$ ) values for the addition at the C–C bonds are different. For the [2+1] addition of an oxygen atom onto  $C_{50}$ , the C<sub>2</sub>–C<sub>3</sub> or C<sub>3</sub>–C<sub>4</sub> bonds with the larger  $K_D$  value is preferable to the C<sub>4</sub>–C<sub>4</sub>' bond, which is consistent with the results calculated by Xu et al.<sup>[8]</sup>

#### Conclusion

The higher reactivities of SWCNTs and fullerene relative to graphite stem from their curvatures. Their variously curved frameworks are formed in two steps, which include the hybridization of the vertex atomic orbitals and bonding between these hybrid orbitals, in which the former is embodied in the mean property of the curvature and the latter in the directional property of the curvature. The directional-curvature theory, developed for the first time herein, covers the mean and directional properties of the curvatures of SWCNTs and fullerenes and allows a unified approach to manifest the strain energies of the vertices and bonds in SWCNTs and fullerenes. The pyramidized angle  $\theta_p$  can not be used to explain the anisotropy of the chemistry of SWCNTs and fullerenes, and the curvature theory contributes to overcoming this difficulty. The theory, in essence, reveals that the  $K_{\rm D}$  and  $K_{\rm M}$  values have an inherent relation to

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the bond curvature K and pyramidized angle  $\theta_p$  and demonstrates that there are approximate curvature conservations in SWCNTs and fullerenes.

Prediction of the binding energies or the corresponding reactivities, which depend directly on the strain energies or  $K_D$  ( $K_M$ ) values, has been verified by a series of theoretical calculations and experiments. It must be pointed out, as the proportional relation is used for the reactivity of the bonds of fullerenes, that consideration of some factors, such as, bonds with [6,6] or [5,6] ring fusion and obviously different lengths, should be required. The surface of many nanoclusters are formed by various orientations of the dangling orbitals relative to the  $\sigma$  framework, thus the directional-curvature theory will be conducive to the study of their reactive mechanisms.<sup>[20]</sup>

#### Acknowledgements

We gratefully acknowledge financial support from the Natural Science Foundation of China (20673019, 20303002), the key project of the Fujian Province (2005HZ01–2–6), and the Natural Educational Department Foundation of PhD Unit (20050386003).

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Received: January 16, 2007 Published online: May 11, 2007

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